

# Numerical Modeling of the Process of Thermoplastic Deformation of Transversally Isotropic Parallelepipeds



M. R. Babajanov, A. A. Kalandarov, U. E. Adambaev

**Abstract:** *The paper proposes a modified version of the iterative method for numerically solving a three-dimensional uncoupled boundary-value problem that describes the process of thermoplastic deformations of a transversely isotropic parallelepiped. A discrete analogue of the boundary value problem is compiled on the basis of the finite-difference method. A recurrent finite-difference relation is written which allows one to find the desired components of the displacement vector in combination with the iterative method. It is assumed that, at a first approximation, the values of the sought displacements in the internal nodes are trivial. The essence of the method is demonstrated by solving the thermoplastic boundary-value problem for a transversely isotropic parallelepiped. The proposed method can be applied to solve related problems of dynamic thermoplasticity.*

**Keywords:** *Coupled problems, displacement, iterative method strain, stress, thermoplasticity.*

## I. INTRODUCTION

Mathematical models describing the process of linear deformation of solids with temperature taken into account, were first considered by Duhamel – Neumann [1]-[3], in which it was assumed that the total deformation consists of elastic deformation and thermal expansion. Questions of the theory of thermoplasticity of deformable solids are investigated in the following papers [4]-[6].

The issues of plastic deformation of solids with regard to the temperature, were first considered in more detail in the works of P.M.Nakhdi[7] and others. Further, these studies were continued in the works of Yu.N. Shevchenko [8]-[10], D. Kolarov [11], J.Casey[12], Aboudi [13],[14], W.F.Chen[15] and others [16],[17].

In many applied engineering and technical problems, the process of plastic deformation of structures and their elements, taking into account temperature influences, can be described by model equations of two types, namely “connected” and “unconnected” thermomechanical problems.

In an unconnected problem, the heat equation responsible for thermal factors is solved separately, and its results are used as a well-known parameter in solving the basic model equations of “thermoplasticity” considered in combination with the equation of motion. Note that in “coupled problems”, in contrast to “uncoupled problems”, a simultaneous solutions of the equations of “thermoplasticity” and the equation of “heat flux” is assumed. Since deformation cause the appearance of temperature, this approach allows one to more adequately describe the process of thermoplastic deformation of structures and their elements under the influence of mechanical and thermal influences. To formulate the above thermoplastic model equations, we need to know the nonlinear relationship between the stress and strain tensors and taking into account temperature [18]-[23].

## II. PROBLEM STATEMENT

Let us consider an unrelated static boundary-value problem describing the process of thermoplastic deformation of a transversely isotropic parallelepiped, based on the deformation theory, which consists of the equilibrium equation

$$\sum_{j=1}^3 \frac{\partial \sigma_{ij}}{\partial x_j} + X_i = 0, \quad i = 1, 2, 3 \quad (1)$$

and, which determines the relation of the deformation theory of transversely isotropic bodies [24], [27], [28]

$$\sigma_{ij} = C_{ijke} \varepsilon_{kl} - \beta_{ij} (T - T_0) \delta_{ij} - 2(\lambda_2 - \lambda_2^*) (1 - \frac{p^*}{p}) p_{ij} - 2(\lambda_5 - \lambda_5^*) (1 - \frac{q^*}{q}) q_{ij} \quad \text{for } p \geq p^*, q \geq q^* \quad (2)$$

Cauchy relations

$$\varepsilon_{kl} = \frac{1}{2} (u_{k,l} + u_{l,k}) \quad (3)$$

and boundary conditions

$$u_i |_{\Sigma_1} = u_i^0, \quad \sum_{j=1}^3 \sigma_{ij} n_j |_{\Sigma_2} = S_i^0 \quad (4)$$

Where  $C_{ijke}$  - tensor of elastic constants of a transversally isotropic material,  $p^*$ ,  $q^*$  - yield strength in the longitudinal and transverse directions of a transversally isotropic body,  $u_i$  - displacement vector components;  $X_i$ ,  $S_i$  - bulk and surface forces, respectively,  $\Sigma_1, \Sigma_2$  - surface parts,  $\Sigma$  - volume  $V$ ,  $n_j$  - external normal to the surface  $\Sigma_2$  - volume  $V$ ,  $T$  -

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absolute temperature,  $\delta_{ij}$ - Kronecker symbol,  $\alpha_{kl}$  - thermal expansion tensor.

$$\begin{aligned} \lambda_1 &= C_{2211}, \lambda_2 = C_{1212}, \lambda_3 = C_{1133}, \lambda_4 = C_{3333}, \lambda_5 = C_{1313}, \\ \lambda_6 &= \lambda_7 = \lambda_1 + 2\lambda_2 = C_{2211} + 2C_{1212}, \\ \lambda_8 &= \lambda_3 = C_{1133}, \lambda_9 = \lambda_5 = C_{1313}. \\ \beta_{ij} &= C_{ijkl}\alpha_{kl}, \quad \alpha_{ij} = \alpha_1(\delta_{ij} - \delta_{i3}\delta_{j3}) + \alpha_3\delta_{i3}\delta_{j3} \end{aligned}$$

Substituting (3) into (2) and obtained in (1), the boundary-value problem (1-4) can be written down with respect to the displacements  $u_i$  ( $u = u_1, v = u_2, w = u_3$ ) in the following form (for simplicity, we neglect volume forces)

$$\left\{ \begin{aligned} &C_{1111} \frac{\partial^2 u}{\partial x^2} + C_{1212} \frac{\partial^2 u}{\partial y^2} + C_{1313} \frac{\partial^2 u}{\partial z^2} + (C_{1122} + C_{1212}) \times \\ &\times \frac{\partial^2 v}{\partial x \partial y} + (C_{1133} + C_{1313}) \frac{\partial^2 w}{\partial x \partial z} - \beta_{11} \frac{\partial T}{\partial x} - F_1 = 0 \\ &C_{1212} \frac{\partial^2 v}{\partial x^2} + C_{2222} \frac{\partial^2 v}{\partial y^2} + C_{2323} \frac{\partial^2 v}{\partial z^2} + (C_{2211} + C_{1212}) \times \\ &\times \frac{\partial^2 u}{\partial x \partial y} + (C_{2233} + C_{2323}) \frac{\partial^2 w}{\partial y \partial z} - \beta_{22} \frac{\partial T}{\partial y} - F_2 = 0 \\ &C_{1313} \frac{\partial^2 w}{\partial x^2} + C_{2323} \frac{\partial^2 w}{\partial y^2} + C_{3333} \frac{\partial^2 w}{\partial z^2} + (C_{3311} + C_{1313}) \times \\ &\times \frac{\partial^2 u}{\partial x \partial z} + (C_{3322} + C_{2323}) \frac{\partial^2 v}{\partial y \partial z} - \beta_{33} \frac{\partial T}{\partial z} - F_3 = 0 \end{aligned} \right. \quad (5)$$

$$\left. \begin{aligned} &u_i|_{\Sigma_1} = u_i^0, \\ &\sum_{j=1}^3 \left\{ \begin{aligned} &\lambda \delta_{ij} \left( \sum_{k=1}^3 \frac{\partial u_k}{\partial x_k} \right) + \mu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \\ &-(3\lambda + 2\mu)\alpha(T - T_0)\delta_{ij} - \sigma_{ij}^p \end{aligned} \right\} n_j \Big|_{\Sigma_2} = S_i^0 \end{aligned} \right. \quad (6)$$

where

$$\sigma_{ij}^p = \begin{cases} 0, & \text{for } p < 0 \text{ and } q < 0, \\ 2(\lambda_2 - \lambda_2')(1 - \frac{p}{p^*})p_{ij}, & \text{for } p \geq p^*, \\ 2(\lambda_5 - \lambda_5')(1 - \frac{q}{q^*})q_{ij}, & \text{for } q \geq q^* \end{cases} \quad (7)$$

$$F_i = \sum_{j=1}^3 \frac{\partial \sigma_{ij}^p}{\partial x_j} 2(\lambda_2 - \lambda_2')(1 - \frac{p}{p^*}) \frac{\partial p_{ij}}{\partial x_j} + 2(\lambda_5 - \lambda_5')(1 - \frac{q}{q^*}) \frac{\partial q_{ij}}{\partial x_j}$$

for  $p \geq p^*, q \geq q^*$

Value  $\sigma_{ij}^p$  represents the nonlinear part of the defining relation (2), and  $F_i$  nonlinear part of equation (1).

### III. NUMERICAL SOLUTION METHOD

Replacing the derivatives in equations (5-6) with the corresponding finite-difference relations, we find the finite-difference analogue [25],[26] of the equilibrium equation (5)

$$\begin{aligned} &C_{1111} \frac{u_{i+1,j,k} - 2u_{i,j,k} + u_{i-1,j,k}}{h_1^2} + \\ &+ C_{1212} \frac{u_{i,j+1,k} - 2u_{i,j,k} + u_{i,j-1,k}}{h_2^2} + \\ &+ C_{1313} \frac{u_{i,j,k+1} - 2u_{i,j,k} + u_{i,j,k-1}}{h_3^2} + \end{aligned} \quad (8)$$

$$\begin{aligned} &(C_{1122} + C_{1212}) \frac{v_{i+1,j+1,k} - v_{i-1,j+1,k} - v_{i+1,j-1,k} + v_{i-1,j-1,k}}{4h_1h_2} + \\ &(C_{1133} + C_{1313}) \frac{w_{i+1,j,k+1} - w_{i-1,j,k+1} - w_{i+1,j,k-1} + w_{i-1,j,k-1}}{4h_1h_3} - \\ &\quad - \frac{\beta_{11}(t_{i+1,j,k} - t_{i-1,j,k})}{2h_1} = F_1 \\ &C_{1212} \frac{v_{i+1,j,k} - 2v_{i,j,k} + v_{i-1,j,k}}{h_1^2} + \\ &+ C_{2222} \frac{v_{i,j+1,k} - 2v_{i,j,k} + v_{i,j-1,k}}{h_2^2} + \\ &+ C_{2323} \frac{v_{i,j,k+1} - 2v_{i,j,k} + v_{i,j,k-1}}{h_3^2} + \end{aligned} \quad (9)$$

$$\begin{aligned} &(C_{2211} + C_{1212}) \frac{u_{i+1,j+1,k} - u_{i-1,j+1,k} - u_{i+1,j-1,k} + u_{i-1,j-1,k}}{4h_1h_2} + \\ &(C_{2233} + C_{2323}) \frac{w_{i,j+1,k+1} - w_{i,j-1,k+1} - w_{i,j+1,k-1} + w_{i,j-1,k-1}}{4h_2h_3} - \\ &\quad - \frac{\beta_{22}(t_{i,j+1,k} - t_{i,j-1,k})}{2h_2} = F_2 \end{aligned}$$

$$\begin{aligned} &C_{1313} \frac{w_{i+1,j,k} - 2w_{i,j,k} + w_{i-1,j,k}}{h_1^2} + \\ &+ C_{2323} \frac{w_{i,j+1,k} - 2w_{i,j,k} + w_{i,j-1,k}}{h_2^2} + \\ &+ C_{3333} \frac{w_{i,j,k+1} - 2w_{i,j,k} + w_{i,j,k-1}}{h_3^2} + \\ &+ (C_{3311} + C_{1313}) \frac{u_{i+1,j,k+1} - u_{i-1,j,k+1} - u_{i+1,j,k-1} + u_{i-1,j,k-1}}{4h_1h_3} + \\ &(C_{3322} + C_{2323}) \frac{v_{i,j+1,k+1} - v_{i,j-1,k+1} - v_{i,j+1,k-1} + v_{i,j-1,k-1}}{4h_1h_2} - \\ &\quad - \frac{\beta_{33}(t_{i,j,k+1} - t_{i,j,k-1})}{2h_3} = F_3 \end{aligned} \quad (10)$$

and boundary conditions, on two opposite faces of the parallelepiped perpendicular to the axis OX

$$\sigma_{11}|_{x=0,l_1} = \left[ \begin{aligned} &C_{1111} \frac{\partial u}{\partial x} + C_{1122} \frac{\partial v}{\partial y} + C_{1133} \frac{\partial w}{\partial z} - \\ &-2(\lambda_2 - \lambda_2')(1 - \frac{p}{p^*})p_{11} - \beta_{11} \frac{\partial T}{\partial x} \end{aligned} \right]_{x=0,l_1} = S_1^0$$

$$\sigma_{12}|_{x=0,l_1} = \left[ \begin{array}{l} C_{1212} \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) - \\ -2(\lambda_2 - \lambda_2^*)(1 - \frac{p^*}{p_u}) p_{12} \end{array} \right]_{x=0,l_1} = S_2^0 \quad (11)$$

$$\sigma_{13}|_{x=0,l_1} = \left[ \begin{array}{l} C_{1313} \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) - \\ -2(\lambda_5 - \lambda_5^*)(1 - \frac{q^*}{q_u}) q_{13} \end{array} \right]_{x=0,l_1} = S_3^0$$

The remaining faces of the parallelepiped are free of loads. The boundary conditions can be set with ease with respect to the displacement vector.

Further, resolving the finite-difference equations (8-10) relatively  $u_{i,j,k}$ ,  $v_{i,j,k}$ ,  $w_{i,j,k}$  accordingly, it is possible to organize the following iterative process:

$$u_{i,j,k}^{(n+1)} = \left[ \frac{C_{1111}}{h_1^2} (u_{i+1,j,k}^{(n)} + u_{i-1,j,k}^{(n)}) + \frac{C_{1212}}{h_2^2} (u_{i,j+1,k}^{(n)} + u_{i,j-1,k}^{(n)}) + \frac{C_{1313}}{h_3^2} (u_{i,j,k+1}^{(n)} + u_{i,j,k-1}^{(n)}) + \frac{C_{1122} + C_{1212}}{4h_1h_2} (v_{i+1,j+1,k}^{(n)} - v_{i+1,j-1,k}^{(n)} - v_{i-1,j+1,k}^{(n)} + v_{i-1,j-1,k}^{(n)}) + \frac{C_{1133} + C_{1313}}{4h_1h_3} (w_{i+1,j,k+1}^{(n)} - w_{i+1,j,k-1}^{(n)} - w_{i-1,j,k+1}^{(n)} + w_{i-1,j,k-1}^{(n)}) - b_{11} \frac{(t_{i+1,j,k} - t_{i-1,j,k})}{2h_1} - F_{1 i,j,k} \right] / \left( \frac{2C_{1111}}{h_1^2} + \frac{2C_{1212}}{h_2^2} + \frac{2C_{1313}}{h_3^2} \right) \quad (12)$$

$$v_{i,j,k}^{(n+1)} = \left[ \frac{C_{1212}}{h_1^2} (v_{i+1,j,k}^{(n)} + v_{i-1,j,k}^{(n)}) + \frac{C_{2222}}{h_2^2} (v_{i,j+1,k}^{(n)} + v_{i,j-1,k}^{(n)}) + \frac{C_{2323}}{h_3^2} (v_{i,j,k+1}^{(n)} + v_{i,j,k-1}^{(n)}) + \frac{C_{2211} + C_{1212}}{4h_1h_2} (u_{i+1,j+1,k}^{(n)} - u_{i+1,j-1,k}^{(n)} - u_{i-1,j+1,k}^{(n)} + u_{i-1,j-1,k}^{(n)}) + \frac{C_{2233} + C_{2323}}{4h_2h_3} (w_{i,j+1,k+1}^{(n)} - w_{i,j-1,k+1}^{(n)} - w_{i,j+1,k-1}^{(n)} + w_{i,j-1,k-1}^{(n)}) - b_{22} \frac{(t_{i,j+1,k} - t_{i,j-1,k})}{2h_2} - F_{2 i,j,k} \right] / \left( \frac{2C_{1212}}{h_1^2} + \frac{2C_{2222}}{h_2^2} + \frac{2C_{1313}}{h_3^2} \right) \quad (13)$$

$$w_{i,j,k}^{(n+1)} = \left[ \frac{C_{1313}}{h_1^2} (w_{i+1,j,k}^{(n)} + w_{i-1,j,k}^{(n)}) + \frac{C_{2323}}{h_2^2} (w_{i,j+1,k}^{(n)} + w_{i,j-1,k}^{(n)}) + \frac{C_{3333}}{h_3^2} (w_{i,j,k+1}^{(n)} + w_{i,j,k-1}^{(n)}) + \frac{C_{3311} + C_{1313}}{4h_1h_3} (u_{i+1,j,k+1}^{(n)} - u_{i+1,j,k-1}^{(n)} - u_{i-1,j,k+1}^{(n)} + u_{i-1,j,k-1}^{(n)}) + \frac{C_{3322} + C_{2323}}{4h_2h_3} (v_{i,j+1,k+1}^{(n)} - v_{i,j-1,k+1}^{(n)} - v_{i,j+1,k-1}^{(n)} + v_{i,j-1,k-1}^{(n)}) - b_{33} \frac{(t_{i,j,k+1} - t_{i,j,k-1})}{2h_3} - F_{3 i,j,k} \right] / \left( \frac{2C_{1313}}{h_1^2} + \frac{2C_{2323}}{h_2^2} + \frac{2C_{3333}}{h_3^2} \right) \quad (14)$$

The values of displacements in the zeroth approximation of the iterative process are considered as a trivial. Recall that, as a first approximation, the terms  $F_{1 i,j,k}$ ,  $F_{2 i,j,k}$ ,  $F_{3 i,j,k}$  are considered zero and the linear problems of the theory of elasticity are solved. Based on the results obtained, the condition for the transition of the elastic limit, i.e.  $p \geq p^*$ ,  $q \geq q^*$  and when they are executed, they are calculated  $F_{1 i,j,k}$ ,  $F_{2 i,j,k}$ ,  $F_{3 i,j,k}$  and, again the linear problem is solved with the new right-hand side, and continues until the required accuracy is achieved. This is the essence of the method of elastic solutions.

#### IV. TEST TASK

When solving the model problem numerically, we considered a transversely isotropic thermoplastic parallelepiped pinched on all sides, inside which a temperature field is given by the following formula:

$$T(x_1, x_2, x_3) = T_0 \sin \frac{\pi x_1}{l_1} \sin \frac{\pi x_2}{l_2} \sin \frac{\pi x_3}{l_3}$$

Elastic modules, hardening modules and elastic limits had the following dimensionless values:

$$C_{1111} = 5.68, C_{1212} = 2.735, C_{1313} = 2.39, C_{1122} = 0.21, C_{1133} = 0.19, C_{2222} = 5.68, C_{2233} = 0.19, C_{2323} = 2.39, C_{3333} = 5.35, \lambda_2^* = 2.5, \lambda_5^* = 2.24, p^* = 0.14, q^* = 0.11, T_0 = 20.$$

Below we consider the propagation of plastic zones in various sections of the transversally isotropic parallelepiped under consideration.

**Table 1. The values of the intensity of the strain tensor  $p$  for  $z=0,7$**

	x=0	x=0.1	x=0.2	x=0.3	x=0.4	x=0.5	x=0.6	x=0.7	x=0.8	x=0.9	x=1
y=0	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
y=0.1	0.00000	0.10830	0.12893	0.14592	0.15313	0.15491	0.15355	0.14648	0.12922	0.10838	0.00000
y=0.2	0.00000	0.12884	0.10128	0.09329	0.09148	0.09124	0.09198	0.09387	0.10153	0.12911	0.00000
y=0.3	0.00000	0.14590	0.09377	0.06079	0.04546	0.04197	0.04561	0.06090	0.09367	0.14604	0.00000
y=0.4	0.00000	0.15299	0.09207	0.04567	0.01788	0.01067	0.01789	0.04561	0.09188	0.15297	0.00000
y=0.5	0.00000	0.15453	0.09153	0.04213	0.01073	0.00002	0.01073	0.04213	0.09153	0.15453	0.00000

y=0.6	0.00000	0.15297	0.09188	0.04561	0.01789	0.01067	0.01788	0.04567	0.09207	0.15299	0.00000
y=0.7	0.00000	0.14604	0.09367	0.06090	0.04561	0.04197	0.04546	0.06079	0.09377	0.14590	0.00000
y=0.8	0.00000	0.12911	0.10153	0.09387	0.09198	0.09124	0.09148	0.09329	0.10128	0.12884	0.00000
y=0.9	0.00000	0.10838	0.12922	0.14648	0.15355	0.15491	0.15313	0.14592	0.12893	0.10830	0.00000
y=1	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

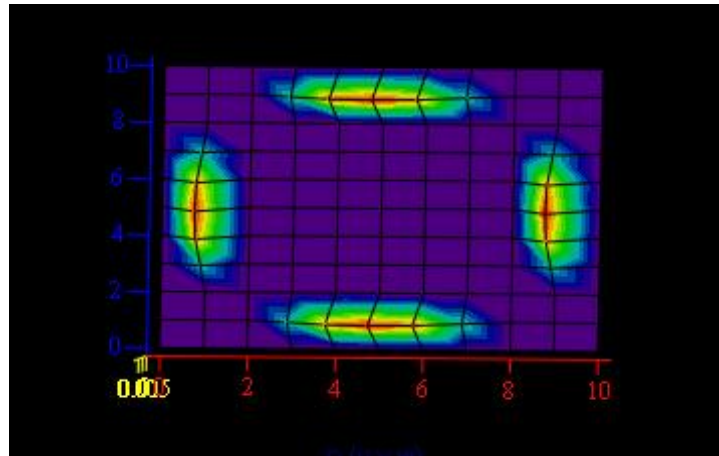


Fig. 1. The zone of plasticity according to the intensity of the tensor deformations  $p$  in the  $XOY$  plane for  $z=0,7$  ( $p \geq p^*$ )

Table 2. The values of the intensity of the strain tensor  $q$  for  $z=0,7$

	y=0	y=0.1	y=0.2	y=0.3	y=0.4	y=0.5	y=0.6	y=0.7	y=0.8	y=0.9	y=1
z=0	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
z=0.1	0.00000	0.05489	0.07683	0.09445	0.10553	0.10926	0.10544	0.09433	0.07678	0.05491	0.00000
z=0.2	0.00000	0.07660	0.08492	0.09123	0.09449	0.09544	0.09449	0.09125	0.08499	0.07688	0.00000
z=0.3	0.00000	0.09410	0.09114	0.08400	0.07547	0.07147	0.07548	0.08404	0.09123	0.09434	0.00000
z=0.4	0.00000	0.10517	0.09436	0.07543	0.05198	0.03842	0.05199	0.07545	0.09441	0.10530	0.00000
z=0.5	0.00000	0.10899	0.09532	0.07142	0.03840	0.00000	0.03840	0.07142	0.09532	0.10899	0.00000
z=0.6	0.00000	0.10530	0.09441	0.07545	0.05199	0.03842	0.05198	0.07543	0.09436	0.10517	0.00000
z=0.7	0.00000	0.09434	0.09123	0.08404	0.07548	0.07147	0.07547	0.08400	0.09114	0.09410	0.00000
z=0.8	0.00000	0.07688	0.08499	0.09125	0.09449	0.09544	0.09449	0.09123	0.08492	0.07660	0.00000
z=0.9	0.00000	0.05491	0.07678	0.09433	0.10544	0.10926	0.10553	0.09445	0.07683	0.05489	0.00000
z=1	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

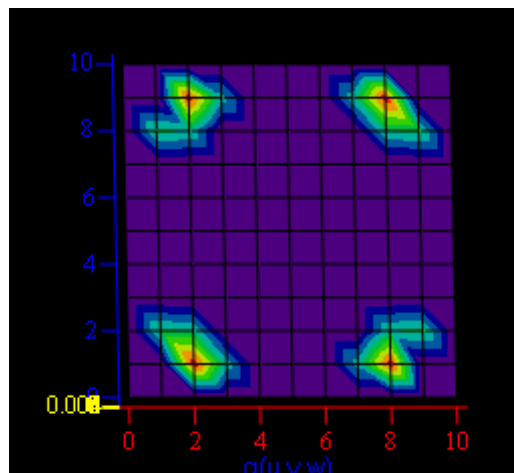
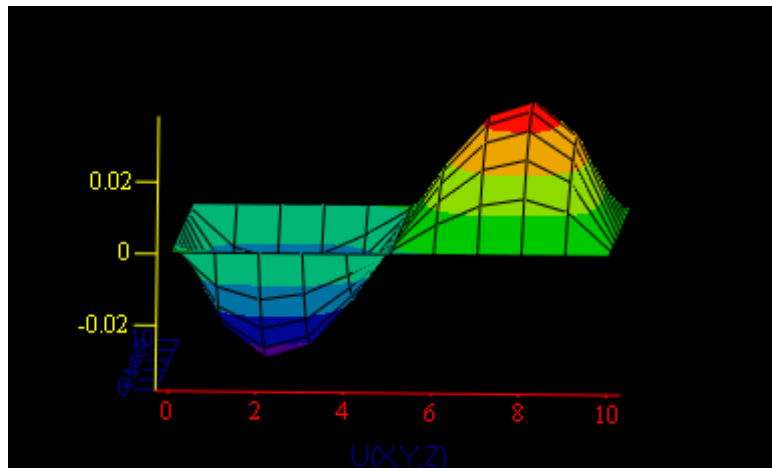


Fig. 2. The zone of plasticity according to the intensity of the tensor deformations  $q$  in the  $XOY$  plane for  $z=0,7$  ( $q \geq q^*$ )

The followings are numerical results and 3D function graphs.  $u(x, y, z)$ ,  $v(x, y, z)$ ,  $w(x, y, z)$ ,  $t(x, y, z)$ ,  $p(x, y, z)$ ,  $q(x, y, z)$ , in the area in question.

**Table 3.**The values of movement  $u(x, y, z)$  for  $z=0,5$

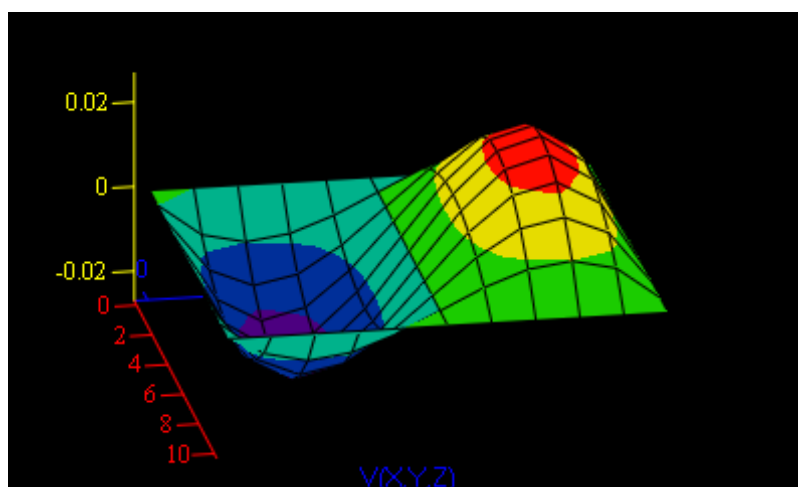
	x=0	x=0.1	x=0.2	x=0.3	x=0.4	x=0.5	x=0.6	x=0.7	x=0.8	x=0.9	x=1
y=0	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
y=0.1	0.00000	-0.01078	-0.01436	-0.01261	-0.00731	0.00000	0.00699	0.01243	0.01435	0.01080	0.00000
y=0.2	0.00000	-0.01768	-0.02385	-0.02110	-0.01216	0.00000	0.01209	0.02107	0.02389	0.01783	0.00000
y=0.3	0.00000	-0.02245	-0.03058	-0.02720	-0.01570	0.00000	0.01569	0.02721	0.03062	0.02257	0.00000
y=0.4	0.00000	-0.02539	-0.03469	-0.03093	-0.01787	0.00000	0.01787	0.03093	0.03472	0.02545	0.00000
y=0.5	0.00000	-0.02641	-0.03609	-0.03219	-0.01861	0.00000	0.01861	0.03219	0.03609	0.02641	0.00000
y=0.6	0.00000	-0.02545	-0.03472	-0.03093	-0.01787	0.00000	0.01787	0.03093	0.03469	0.02539	0.00000
y=0.7	0.00000	-0.02257	-0.03062	-0.02721	-0.01569	0.00000	0.01570	0.02720	0.03058	0.02245	0.00000
y=0.8	0.00000	-0.01783	-0.02389	-0.02107	-0.01209	0.00000	0.01216	0.02110	0.02385	0.01768	0.00000
y=0.9	0.00000	-0.01080	-0.01435	-0.01243	-0.00699	0.00000	0.00731	0.01261	0.01436	0.01078	0.00000
y=1	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000



**Fig. 3:** Distribution graph of the function  $u(x, y, z)$  in the plane XOY for  $z=0,5$

**Table 4.** Values of displacement  $v(x, y, z)$  for  $z=0,8$

	x=0	x=0.1	x=0.2	x=0.3	x=0.4	x=0.5	x=0.6	x=0.7	x=0.8	x=0.9	x=1
y=0	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
y=0.1	0.00000	-0.00762	-0.01250	-0.01595	-0.01806	-0.01878	-0.01806	-0.01595	-0.01250	-0.00762	0.00000
y=0.2	0.00000	-0.00992	-0.01674	-0.02155	-0.02448	-0.02547	-0.02448	-0.02155	-0.01674	-0.00992	0.00000
y=0.3	0.00000	-0.00859	-0.01470	-0.01901	-0.02163	-0.02252	-0.02163	-0.01901	-0.01470	-0.00859	0.00000
y=0.4	0.00000	-0.00488	-0.00842	-0.01091	-0.01243	-0.01295	-0.01243	-0.01091	-0.00842	-0.00488	0.00000
y=0.5	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
y=0.6	0.00000	0.00488	0.00842	0.01091	0.01243	0.01295	0.01243	0.01091	0.00842	0.00488	0.00000
y=0.7	0.00000	0.00859	0.01470	0.01901	0.02163	0.02252	0.02163	0.01901	0.01470	0.00859	0.00000
y=0.8	0.00000	0.00992	0.01674	0.02155	0.02448	0.02547	0.02448	0.02155	0.01674	0.00992	0.00000
y=0.9	0.00000	0.00762	0.01250	0.01595	0.01806	0.01878	0.01806	0.01595	0.01250	0.00762	0.00000
y=1	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000



**Fig. 4.**Graph of the distribution of the function  $v(x, y, z)$  in the plane XOY for  $z=0,8$

Table 5. The displacement values  $w(x, y, z)$  for  $z=0,3$

	x=0	x=0.1	x=0.2	x=0.3	x=0.4	x=0.5	x=0.6	x=0.7	x=0.8	x=0.9	x=1
y=0	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
y=0.1	0.00000	-0.00714	-0.01273	-0.01690	-0.01950	-0.02038	-0.01949	-0.01688	-0.01272	-0.00713	0.00000
y=0.2	0.00000	-0.01272	-0.02284	-0.03043	-0.03517	-0.03679	-0.03517	-0.03042	-0.02283	-0.01272	0.00000
y=0.3	0.00000	-0.01689	-0.03043	-0.04055	-0.04689	-0.04905	-0.04689	-0.04055	-0.03043	-0.01689	0.00000
y=0.4	0.00000	-0.01950	-0.03518	-0.04689	-0.05424	-0.05674	-0.05423	-0.04689	-0.03517	-0.01950	0.00000
y=0.5	0.00000	-0.02039	-0.03679	-0.04905	-0.05674	-0.05937	-0.05674	-0.04905	-0.03679	-0.02039	0.00000
y=0.6	0.00000	-0.01950	-0.03517	-0.04689	-0.05423	-0.05674	-0.05424	-0.04689	-0.03518	-0.01950	0.00000
y=0.7	0.00000	-0.01689	-0.03043	-0.04055	-0.04689	-0.04905	-0.04689	-0.04055	-0.03043	-0.01689	0.00000
y=0.8	0.00000	-0.01272	-0.02283	-0.03042	-0.03517	-0.03679	-0.03517	-0.03043	-0.02284	-0.01272	0.00000
y=0.9	0.00000	-0.00713	-0.01272	-0.01688	-0.01949	-0.02038	-0.01950	-0.01690	-0.01273	-0.00714	0.00000
y=1	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

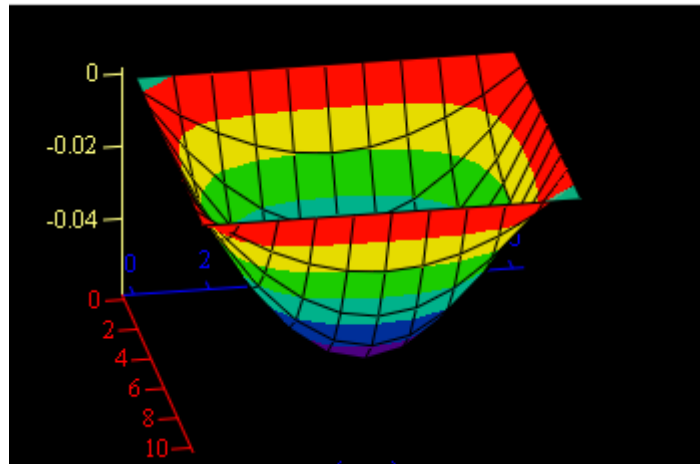


Fig. 5. Distribution graph of the function  $w(x, y, z)$  in the plane XOY for  $z=0,3$

Table 6. Temperature  $t(x, y, z)$  for  $z=0,4$

	x=0	x=0.1	x=0.2	x=0.3	x=0.4	x=0.5	x=0.6	x=0.7	x=0.8	x=0.9	x=1
y=0	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
y=0.1	0.00000	1.81636	3.45492	4.75528	5.59017	5.87785	5.59017	4.75528	3.45492	1.81636	0.00000
y=0.2	0.00000	3.45492	6.57164	9.04508	10.63314	11.18034	10.63314	9.04508	6.57164	3.45492	0.00000
y=0.3	0.00000	4.75528	9.04508	12.44949	14.63525	15.38842	14.63525	12.44949	9.04508	4.75528	0.00000
y=0.4	0.00000	5.59017	10.63314	14.63525	17.20477	18.09017	17.20477	14.63525	10.63314	5.59017	0.00000
y=0.5	0.00000	5.87785	11.18034	15.38842	18.09017	19.02113	18.09017	15.38842	11.18034	5.87785	0.00000
y=0.6	0.00000	5.59017	10.63314	14.63525	17.20477	18.09017	17.20477	14.63525	10.63314	5.59017	0.00000
y=0.7	0.00000	4.75528	9.04508	12.44949	14.63525	15.38842	14.63525	12.44949	9.04508	4.75528	0.00000
y=0.8	0.00000	3.45492	6.57164	9.04508	10.63314	11.18034	10.63314	9.04508	6.57164	3.45492	0.00000
y=0.9	0.00000	1.81636	3.45492	4.75528	5.59017	5.87785	5.59017	4.75528	3.45492	1.81636	0.00000
y=1	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

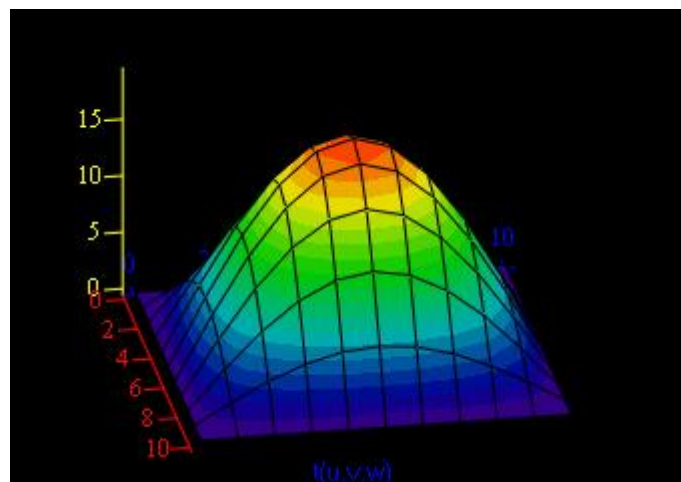


Fig. 6. The distribution graph of the function  $t(x, y, z)$  in the plane XOY for  $z=0,4$

Since the boundary conditions are given symmetrically, the numerical results given in tables (1-6) and figures (1-6) obtained by the above iterative method are

distributed symmetrically in the parallelepiped under consideration,

which ensures the validity of the numerical results and the proposed iterative approach for the numerical solution of thermoplastic problems for transversely isotropic bodies.

## V. CONCLUSION

A three-dimensional thermoplastic boundary value problem for a transversely isotropic parallelepiped is formulated. A discrete analogue of the problem is compiled by the finite-difference method. Finite-difference equations are resolved with respect to the main nodal displacements and a recurrence relation is obtained that is solved by the iterative method in combination with the elastic solution method. The thermoplastic boundary-value problem of a clamped transversally isotropic parallelepiped is numerically solved.

3D graphs of the distribution of the displacement functions  $u(x, y, z)$ ,  $v(x, y, z)$ ,  $w(x, y, z)$  and temperature  $T(x, y, z)$  are constructed. The propagation of plasticity zones in various sections of the transversally isotropic parallelepiped under the influence of a temperature field is studied.

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2. Numerical solution of a Three-Dimensional Coupled thermoplastic problem based on deformation theory. The international journal of science & technology. Vol 7 Issue 9, DOI:10.24940/thejst/2019/v7/ i9/ST1909-012 September, 2019



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